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SCRAM: A Fast Computational Model for the Optical Performance of Point Focus Solar Central Receiver Systems

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SCRAM: A Fast Computational Model for the
Optical Performance of Point Focus
Solar Central Receiver Systems

by

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ABSTRACT

Because of the complexities of heliostat shadowing and blocking calculations, computational models for the optical performance of point focus central receiver (PFCR) systems tend to be too slow for many important applications, such as optimization studies based on performance with realistic weather data. In this paper, a mathematical approximation procedure, designated Sandia Central Receiver Approximation Model (SCRAM) will be described. Rather than simulating the system components from first principles, it relies on data generated by the DELSOL code of Dellin and Fish for the optical performance of PFCR systems, and abstracts a mathematical model using a stepwise regression procedure. The result is a computational procedure which allows the user to define the heliostat field boundaries and tower height arbitrarily, generating a model for optical field performance, including shadowing, blocking, cosine losses, and atmospheric attenuation, and which requires only a polynomial evaluation for each set of sun angles. A comparison with DELSOL for three different fields on three representative days indicates that the rms error of the approximation is 1-3% and that the new code is 1,000-3,000 times as fast as DELSOL. It is also shown that one reason that the accuracy in field performance predictions is higher than that of the generating function for the model is that much of the error in the generating function is due to an oscillatory behavior associated with a moire' pattern in the optical response of the heliostat field.

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I. Introduction

Point focus central receiver (PFCR) systems, consisting of a field of individual heliostats which reflect and concentrate sunlight onto a heat exchanger at the top of a tower, are being seriously considered for large scale production of electricity and/or heat from solar energy. Though this "power tower" concept is simple, the mathematical modeling of optical performance is considerably more difficult than other sun-tracking systems (e.g. parabolic troughs or dishes). Primarily this is because the effects of tower blocking and sun shading depend on the precise orientation of each heliostat with respect to its neighbors, the tower, and the sun, so that complex computer codes are required to treat the problem realistically. Several such codes now exist,^{1,2,3} but they tend to require large computers and long running times. For some purposes (e.g. simulations in which hourly weather data is sampled to determine annual system performance at a particular location) the long running time of even the fastest of these codes³ is prohibitive. The purpose of this paper is to describe a method for constructing an approximate mathematical model for the optical performance of a PFCR which requires very little computation time.

The approach is to make use of the existing simulation codes to generate a large amount of data which can be pictured as a surface in a many-parameter space. We then find a polynomial function which is a good approximation to this surface over the range of parameters of interest. The resulting model is called SCRAM (Sandia Central Receiver Approximation Model) and it is fast enough and accurate enough (compared to the original data base) that it can be used in simulations which sample weather and load variations and in system optimization studies which require thousands or hundreds of thousands of calls to the computational PFCR performance model in each run. An additional benefit is that SCRAM can be used for simple calculations on any computer regardless of size or speed.

It is clear that such an approximation model can be no more accurate or reliable than the simulation code which generates the data base. We will rely primarily on the DELSOL code developed for large applications

by Dellin and Fish,³ and the resulting approximation model can be considered to be an adjunct to that code.

II. Problem Definition

There are a large number of parameters which affect the performance of a PFCR system. These include the position of the sun and its shape, the attenuation characteristics of the atmosphere, the layout of the mirrors on the heliostats and of the heliostats on the field, the tower height, the receiver dimensions and thermal characteristics, the reflectivities of the mirrors and receiver, the tracking and focusing accuracies, and many more. For the sake of discussion, let us say there are N such parameters. The goal of our calculation is to find an approximate polynomial expression for the energy collected as a function of these N parameters, but it is impractical to include all the parameters as independent variables in the polynomial. For one thing, the time and expense required to generate the data base via the large simulation code would be excessive (a reasonable sampling of the N -dimensional space is required). For another, the process of searching for a good polynomial fit to the data would be very difficult. And finally, the resulting model would be too cumbersome for many of the purposes envisioned.

Thus, it is necessary to select a subset of n parameters to be treated as independent variables, and to choose values for the other $(N-n)$ parameters which will be held constant. This selection process is a matter of judgment and depends on the use for which the model is intended. For our model the following simplifications will be made:

- a) A particular heliostat design will be chosen for the model. Any variation in the heliostat parameters requires the generation of a new data base and the construction of a new model.
- b) Similarly, a particular heliostat layout pattern is chosen for the model. The boundaries of the field will be variable

but the heliostat separations at each point in the field are not. Details of the parameters chosen for heliostat design and layout for the model which will be generated in Section III are given in Appendix A.

The efficiency, η , of a PFCR system is defined as the power absorbed by the receiver heat transfer fluid divided by the solar power incident on an area equal to the mirror area oriented normal to the sun's rays. We can write this as a product of the field efficiency and the receiver efficiency and break each of these up into terms associated with different power loss mechanisms as follows:

$$\eta = \eta_{\text{field}} \cdot \eta_{\text{receiver}} \quad (1)$$

$$\eta_{\text{field}} = \eta_1 \cdot \eta_2 \cdot \eta_3 \cdot \eta_4 \cdot \eta_5 \quad (2)$$

$$\eta_{\text{receiver}} = \eta_6 \cdot \eta_7 \cdot \eta_8 \cdot \eta_9 \cdot \eta_{10} \quad (3)$$

Table I indicates the loss mechanism associated with each term.

Table I.

Term	Loss Mechanism
η_1	Cosine--mirror is not normal to sun's rays
η_2	Shadowing--shadow of one heliostat falls on another
η_3	Blocking--light reflected from one heliostat is intercepted by another
η_4	Absorption--at mirror surface
η_5	Attenuation--in the air between heliostat and receiver
=====	
η_6	Spillage--reflected light misses receiver aperture
η_7	Reflection--from receiver surface
η_8	Radiation--due to elevated temperature of receiver surface
η_9	Convection--loss to air adjacent to receiver
η_{10}	Conduction--loss to receiver support structure

In the present analysis, we will achieve an additional simplification by modeling only η_{field} . There are several reasons for this:

- a) Substantial mathematical simplification is achieved.
- b) The dominant losses are the field losses.
- c) Thermal loss mechanisms ($\eta_8, \eta_9, \eta_{10}$) are poorly understood in the relevant parameter regimes, and there are, at this time, no reliable models for these processes (particularly for convective losses).⁴

The intention is not to ignore receiver losses, but rather to separate the problem into two parts and deal only with that part which existing codes (e.g. DELSOL) treat well. Our model for η_{field} can then be coupled with realistic models for the receiver and/or the load. For illustrative purposes, a simplified receiver model will be introduced in Section V. and described in Appendix C. All terms in η_{field} scale geometrically with tower height, except η_5 , which will be treated separately. This means that we can use the tower height, h , as the unit of length throughout the calculation until the final evaluation of total power incident on the receiver region.

We have thus reduced the problem to two key mathematical steps. First, we generate (using DELSOL) a set of data points describing the surface

$$\eta_s(r, \beta, \theta, \phi) \equiv \eta_1 \cdot \eta_2 \cdot \eta_3 \cdot \eta_4 \quad (4)$$

which is the local field efficiency at each radius (r) and azimuth (β) in the field for a particular set of sun polar (θ) and azimuthal (ϕ) angles. (The tower is located at $r = 0$ --see Figure 1.) Second, we find a polynomial approximation for this surface.

$$\eta_p(r, \beta, \theta, \phi) \equiv \eta_s(r, \beta, \theta, \phi) \quad (5)$$

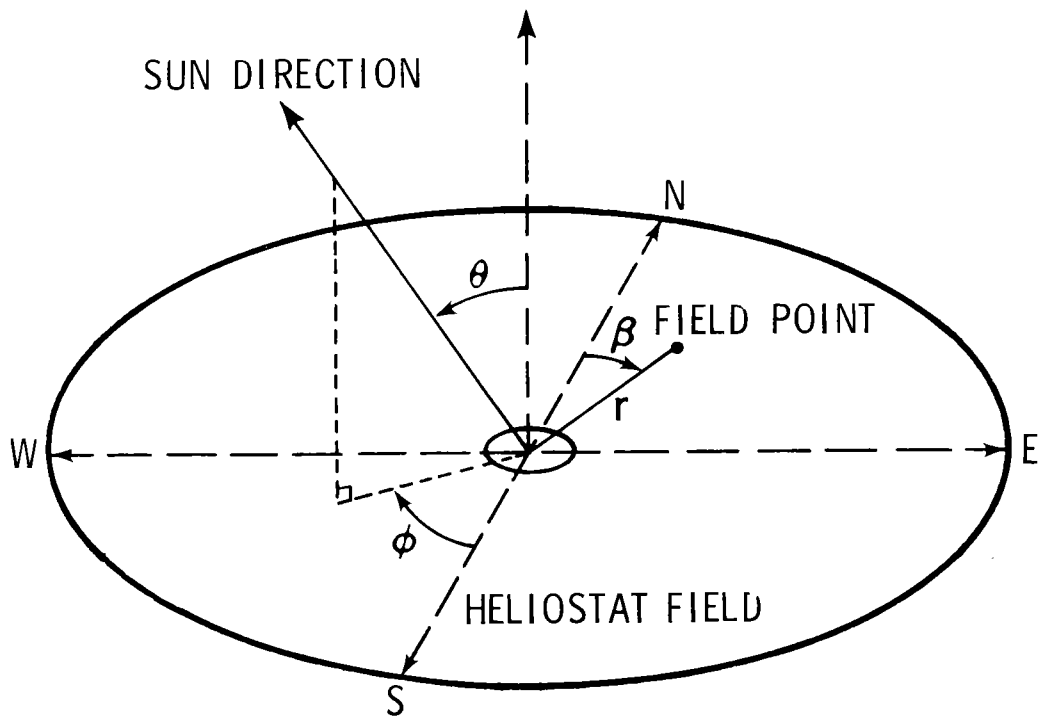


Figure 1. Field and sun angles for analysis

With this we can calculate the field efficiency for any field boundaries:

$$\eta_f(\theta, \phi) = C_f(\theta, \phi)/A_m$$

where A_m is the total mirror area (in m^2) and

$$C_f(\theta, \phi) = h^2 \int \int_{\text{field}} \eta_5(r, h) \rho(r, \beta) \eta_p(r, \beta, \theta, \phi) r dr d\beta \quad (6)$$

where $\rho(r, \beta)$ is the field mirror density (ratio of mirror area to land area). If D is the direct normal insolation (watt/m^2), then the power reflected from the field incident on the receiver area is simply

$$P_f = C_f D$$

The total power, P , is obtained by subtracting from P_f the losses associated with the terms η_6 through η_{10} .

III. Generation of the SCRAM model

There are many ways to generate an approximate fit to a function of several variables. The procedure described below was chosen because of its suitability to the behavior of η_s and for easy implementation of the integration in Eq. (6). We first note that symmetry considerations reduce the number of independent variables to three: r , θ , and $\xi = \phi - \beta$. We then assume

$$\eta_p = \sum_{i,j,k} C_{ijk} U_i(\theta) V_j(\xi) W_k(r) \quad (7)$$

where U , V , and W are the base functions for the approximation optimization. The indices i, j, k will range from 1 to 4 and the choice of base functions is shown in Table II. We have made use of the symmetry of η_s about $\xi = 0$ by using only even functions of ξ in the sum (7). (The odd functions U_5 , U_6 , and U_7 as defined in Table II will be used in subsequent calculations.)

Table II.

Index	$U_i(\xi)$	$U_{i+4}(\xi)$	$V_j(\theta)$	$W_k(r)$
1	1	$\sin \xi$	1	r^{-1}
2	$\cos(\xi)$	$\sin(2\xi)$	$\cos \theta$	1
3	$\cos(2\xi)$	$\sin(3\xi)$	$\sin \theta$	r
4	$\cos(3\xi)$	-	$\cos\theta\sin\theta$	r^2

There are 64 terms in the sum in Eq. (7). The goal is to find a set of coefficients C_{ijk} , most of which are zero, yet which give an acceptable fit to the surface η_s . For this we use a stepwise regression program developed at Kansas State University and Sandia Laboratories.⁵ One can think of the procedure as mapping η_s onto the 64-dimensional hyperspace defined by the variables specified in Table II, and then finding the best fit (in the least squares sense) for any M-dimensional hyperplane as M is sequentially increased from 1 to 2 to 3, etc. It is generally found that there is a point of diminishing returns around $M = 10$, after which relatively little reduction in the root mean square (rms) error of the approximation occurs.

This procedure has been applied to data generated by DELSOL to construct a model for a particular heliostat and layout pattern (detailed parameters are given in Appendix A). Table III contains the coefficients C_{ijk} for the resulting model. All coefficients not listed are zero. (It should be noted that this model is restricted in validity to the range $.4 < r < 8.0$ and its use beyond this range would be an extrapolation from the original DELSOL data.) The rms error of the approximation (compared to the original data base) is .041, which corresponds to a relative error of 6.5%. For reasons which we will discuss later, the relative error in the predictions for field performance will be considerably less than this error in local efficiency.

Table III.

i	j	k	C_{ijk}
1	1	1	-.9260
1	1	2	.5466
2	1	1	.9343
2	1	2	.1114
3	1	1	.5656
3	2	3	.0824
3	3	2	-.02308
4	2	2	.07426
3	2	4	-.005781

The next step is to develop a procedure which will integrate the local efficiency η_p over the field for some specified field boundaries. It is desirable that this integration be done without specifying θ and ϕ so that once a field has been set up, multiple calls to the performance model can be made which require only a polynomial evaluation, not a numerical integration. For this reason we rewrite Eq. (6) in the following way.

$$C_f(\theta, \phi) = \sum_{\substack{i=1,4 \\ j=1,7}} X_{ij} V_i(\theta) U_j(\phi) \quad (8)$$

where the coefficients X_{ij} are sums of integrals over the field variables, r and β , and are independent of θ and ϕ . Expressions for the X_{ij} and a derivation of Eq. (8) are given in Appendix B. Once these coefficients have been evaluated for a particular field configuration, the optical performance of the field for any sun angles requires only the evaluation of the polynomial in Eq. (8), which requires very little computer time.

IV. Moire' Patterns in Optical Performance

Before the details of the procedure for implementing the analysis of the preceding section are given, it is relevant to point out an unusual characteristic of the optical properties of a PFCR system. In Figure 2 we have plotted the local heliostat field efficiency $\eta_s(r, \beta, \theta, \phi)$ (see Eq. (4)) vs. $\xi = \phi - \beta$ for a fixed r and θ , as predicted by DELSOL. Superimposed on the slow (long wavelength) variation, which we associate with cosine losses, is a short wavelength oscillation having 6 cycles from 0° to 360° .

All of this variation can be found to reside in the product $\eta_2 \cdot \eta_3$, due to heliostat shadowing and blocking. It can be better understood by referring to Figure 3, where $(1 - \eta_2 \eta_3)$ is plotted as the dot density in a polar plot of r and ξ for a fixed θ . (A large value of θ , corresponding to a time near sunset, was chosen so that shadowing effects are emphasized.) Fig. 3 represents the actual visual appearance of the heliostat field at a particular moment in time, and what we see is a moire' pattern, which typically results from the superposition of two patterns which are spatially periodic with approximately the same wavelength but phase shifts which result in dark and light bands (corresponding to alternating regions of destructive and constructive interference). In this case the two superimposed patterns are the heliostats and their shadows.

Each lobe of the moire' pattern corresponds to a different mode of shadowing. For example, the horizontal lobe corresponds to the shadowing of one heliostat by a neighbor in the same ring, while the one at $\xi = 45^\circ$ is due to a shadow cast from one ring to another on a diagonal from the radius vector. (Since the azimuthal heliostat separation is constant with radius, while the radial separation increases with increasing radius, all the lobes except the horizontal one disappear beyond a certain radius.⁶)

This behavior has important implications for an approximation procedure of the kind under consideration here. In order to follow the short wavelength variation, it would be necessary to include a large number of short wavelength base functions in the sum in Eq. (7); the four functions

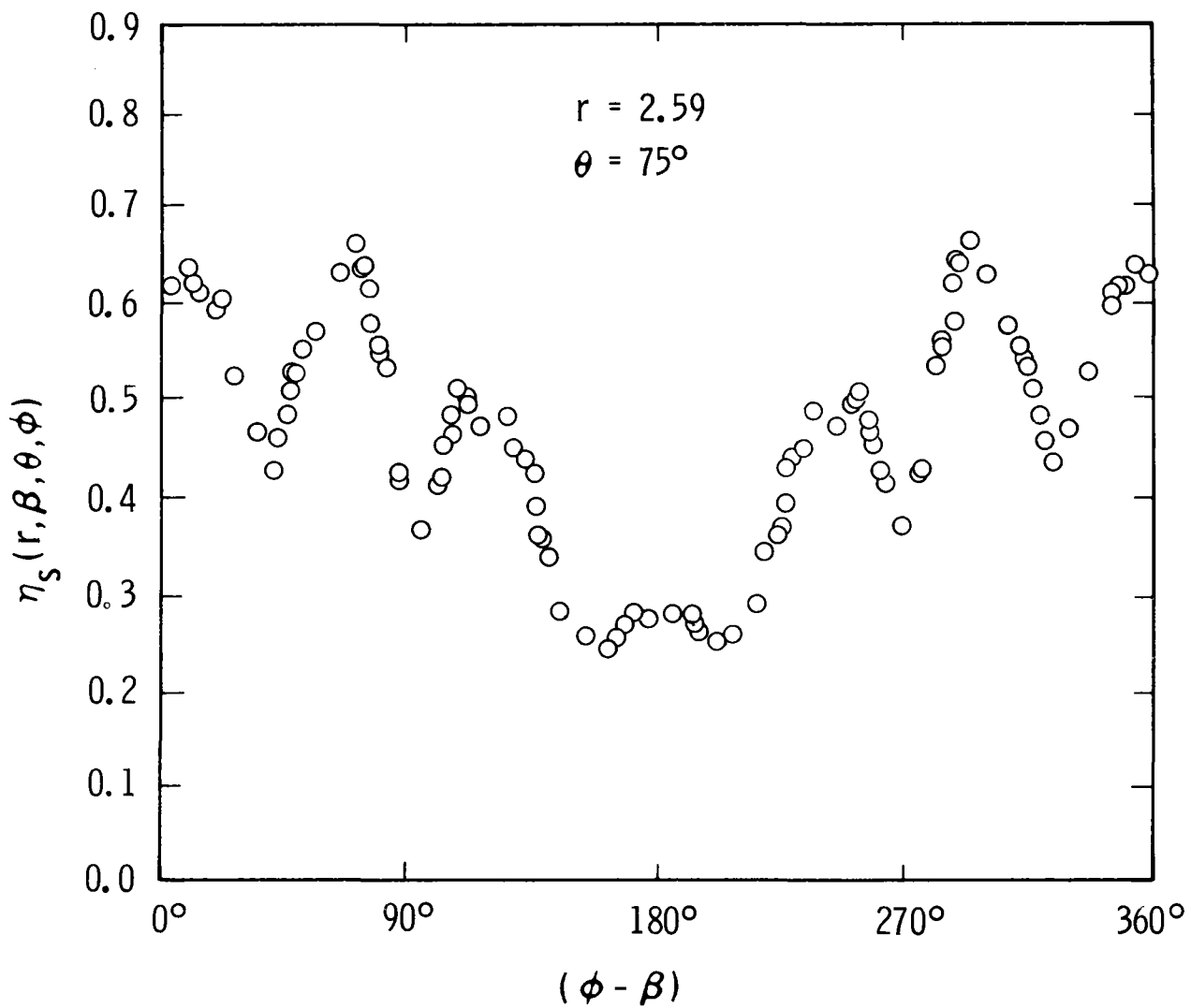


Figure 2. Local field efficiency as predicted by DELSOL, as a function of $(\phi - \beta)$ at a fixed radius in the field, and for a fixed zenith angle.

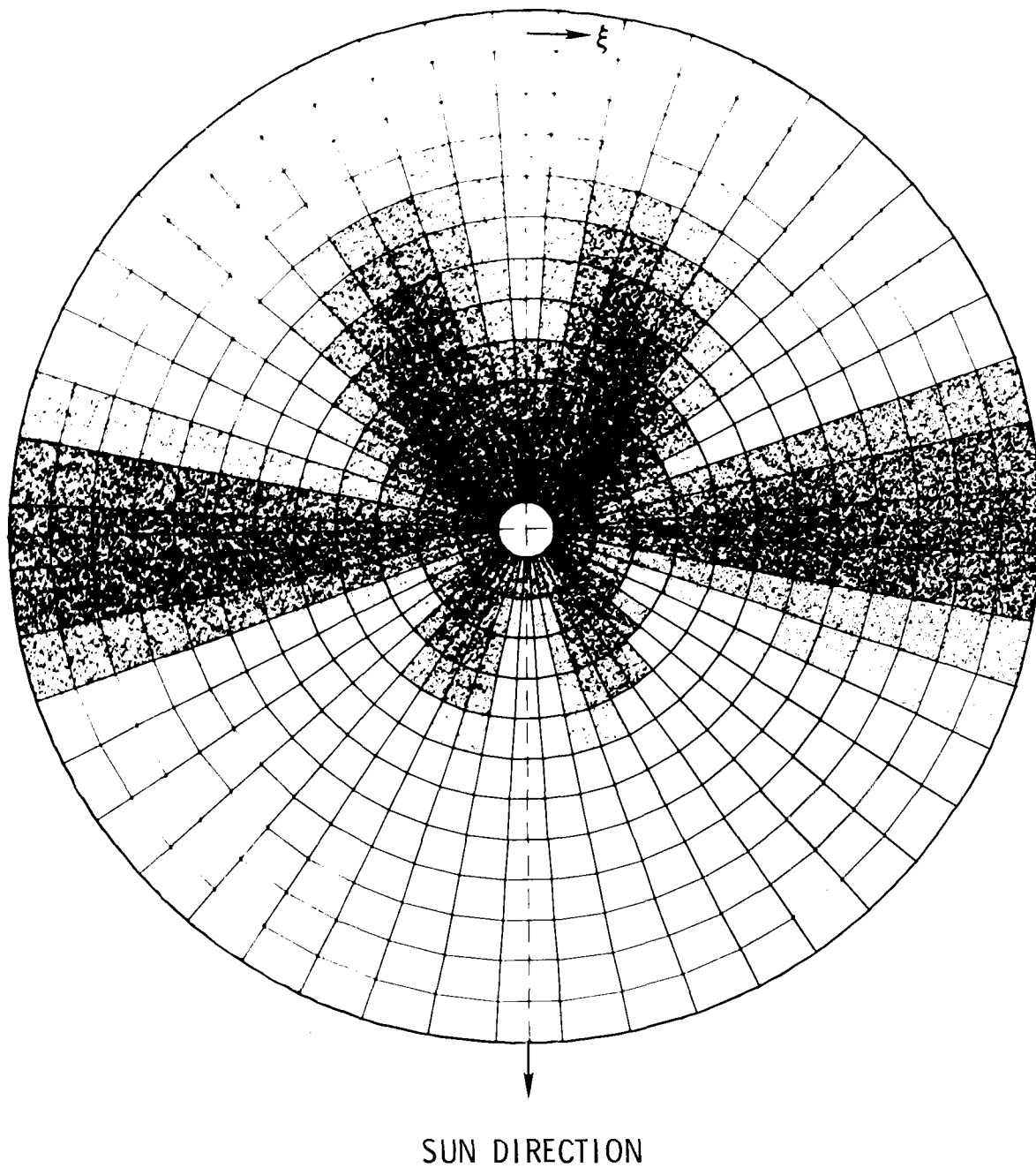


Figure 3. Blocking and shadowing losses plotted as the dot density on a polar plot of the field variables (r, β) . This moiré pattern results from the alignment of the shadow pattern with the mirror pattern.

$U_1(\xi)$ would not be nearly adequate. If this were done, the entire procedure would become considerably more complex and cumbersome. If it were not done, the approximating function would be a smoothed version of the curve shown in Figure 2. The short wavelength oscillation would then contribute significantly to the rms error of the approximation for η_s .

However, because this deviation is oscillatory, its contribution to the error in the approximation for the total field performance is expected to be considerably less than if it were due to random deviation. This is because in the integration over field angle the effect of a positive deviation is effectively cancelled by the negative deviation in the next half-cycle of the oscillation. We will see in the next section that the relative error in the approximation for field performance is considerably less than that for local efficiency.

V. Examples and Tests of the Approximation Model

There are numerous ways that the procedure described in Section III could be implemented to construct a fast performance model for PFCR systems. To demonstrate the use of the procedure and to test its accuracy, we have written a program which allows the specification of arbitrary field boundaries (within the limits of $.4 < r < 8.$) and which samples hourly weather data from the TMY tapes⁷ to determine annual performance characteristics of the field. The logical structure of the program is shown in Figure 4. All of the PFCR modeling functions are contained in the subroutines FLDSET, DENS, and SCRAM so that they could be used unaltered for a number of different purposes with different calling programs. Appendix C contains a listing of these programs, (though it should be noted that for applications on systems other than the one at Sandia Laboratories the subroutines QNC3 and AVINT may have to be replaced by other numerical quadrature routines).

The philosophy behind our analysis is that it is not necessary to model receiver performance at the same time as field performance. It is assumed to be adequate for system design studies, for example, to optimize the field on the basis of field efficiency, and subsequently to optimize receiver

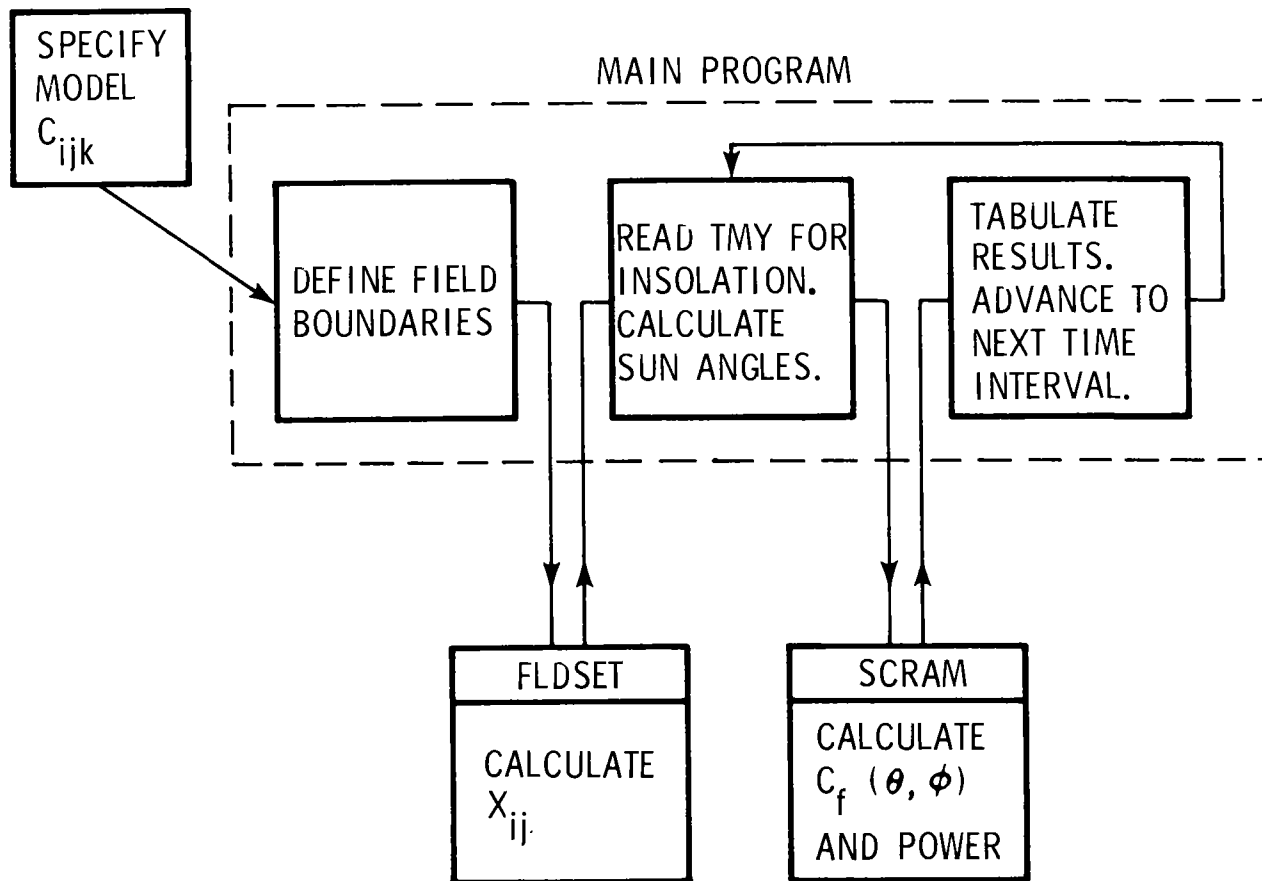


Figure 4. Logical structure of the program which uses the SCRAM model to predict annual performance of PFCR systems for realistic weather data.

design on the basis of the fixed field design. The reason for this effective decoupling is that the dominant cost of the PFCR is the heliostat field. To give an example of how a receiver model would be incorporated in this analysis, we have included a simple example in the program. Receiver losses are written as the sum of two terms: one proportional to the insolation (which includes η_6 and η_7) and the other a constant (including η_8 , η_9 , and η_{10} , which are fixed for a given receiver operating temperature). Details are presented in Appendix C.

We have tested the accuracy of our method by evaluating the performance of three different fields on three different days of the year, and comparing with the predictions of DELSOL for the same fields and days. A clear sky weather model was employed. The results are shown in Figures 1, 2, and 3. As can be seen, the approximation is very good over the entire range of parameters. In fact, it is a "better" approximation, in the least squares sense, than the original approximation for the local efficiency $\eta_s(r, \beta, \theta, \phi)$. We can see this by calculating the relative rms error, which we define as

$$e = \left\{ N^{-1} \sum_{i=1}^N (1 - P_{si}/P_{Di})^2 \right\}^{1/2} \quad (9)$$

where P_{si} is the SCRAM prediction for power collected at the i 'th time interval, and P_{Di} is the corresponding DELSOL prediction.

As indicated in Figures 5-7, e is 3% or less for all cases considered, which should be compared with the 6.5% error characterizing the local field efficiency approximation (see Section III.). The reason for the reduction in the error is that not all of the deviation of η_p from η_s is random-- a significant portion comes from the shadowing and blocking moiré pattern, as discussed in Section IV. The level of 3% error is probably more than adequate for most PFCR modeling requirements, since the uncertainty in many of the parameters used as input to DELSOL are greater. (For example, the variation in heliostat reflectivity due to accumulation of dust and dirt between washings will be considerably higher.⁸)

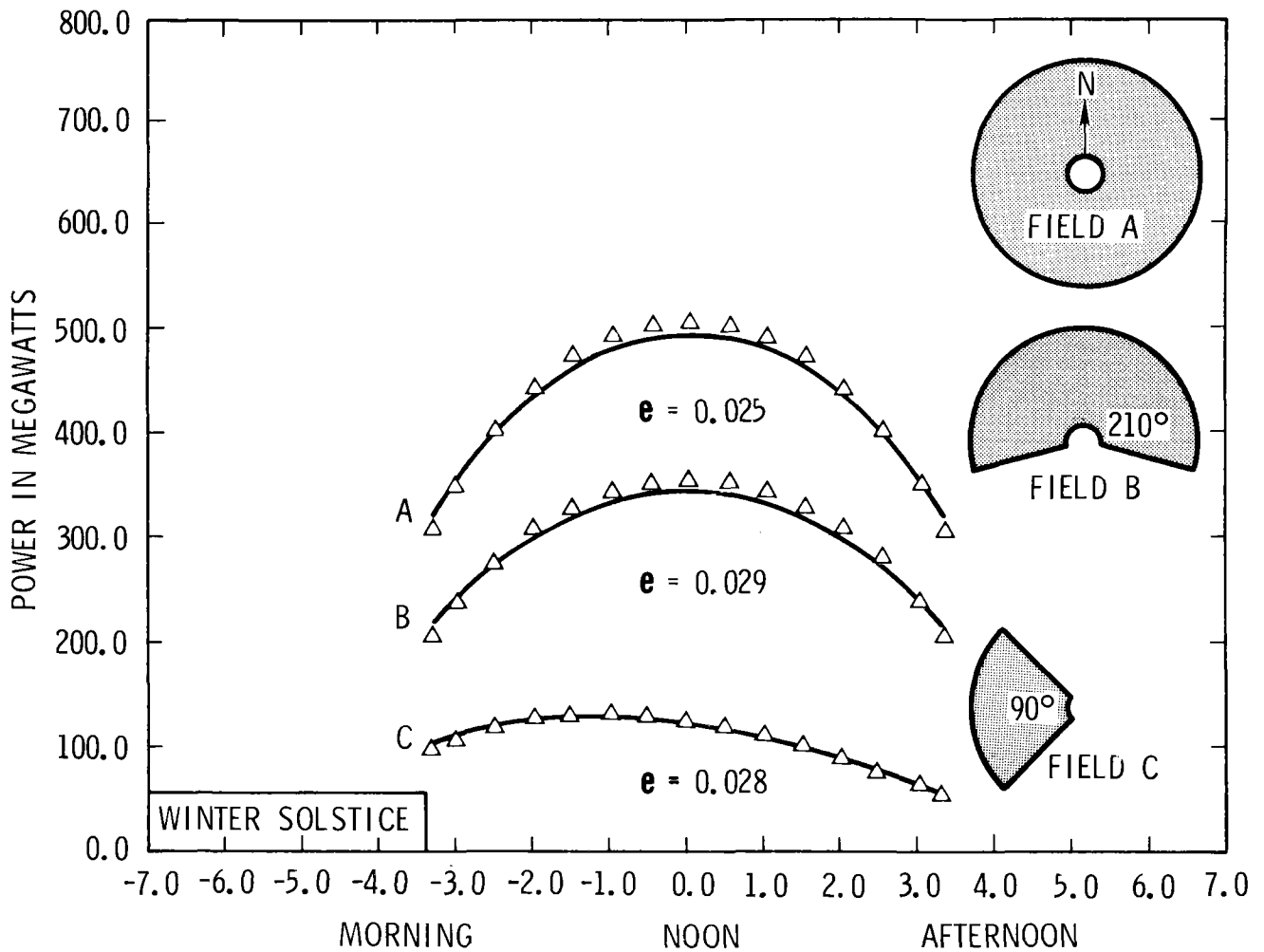


Figure 5. Power delivered to receiver area on December 22 (clear day) for three different field configurations. In all three fields the inner and outer radii are .75 and 7.5 tower heights, and the tower height is 87 m. Latitude is 35° N. The triangles are DELSOL results and the solid line is the SCRAM approximation. e is the rms error of the approximation, relative to the mean value.

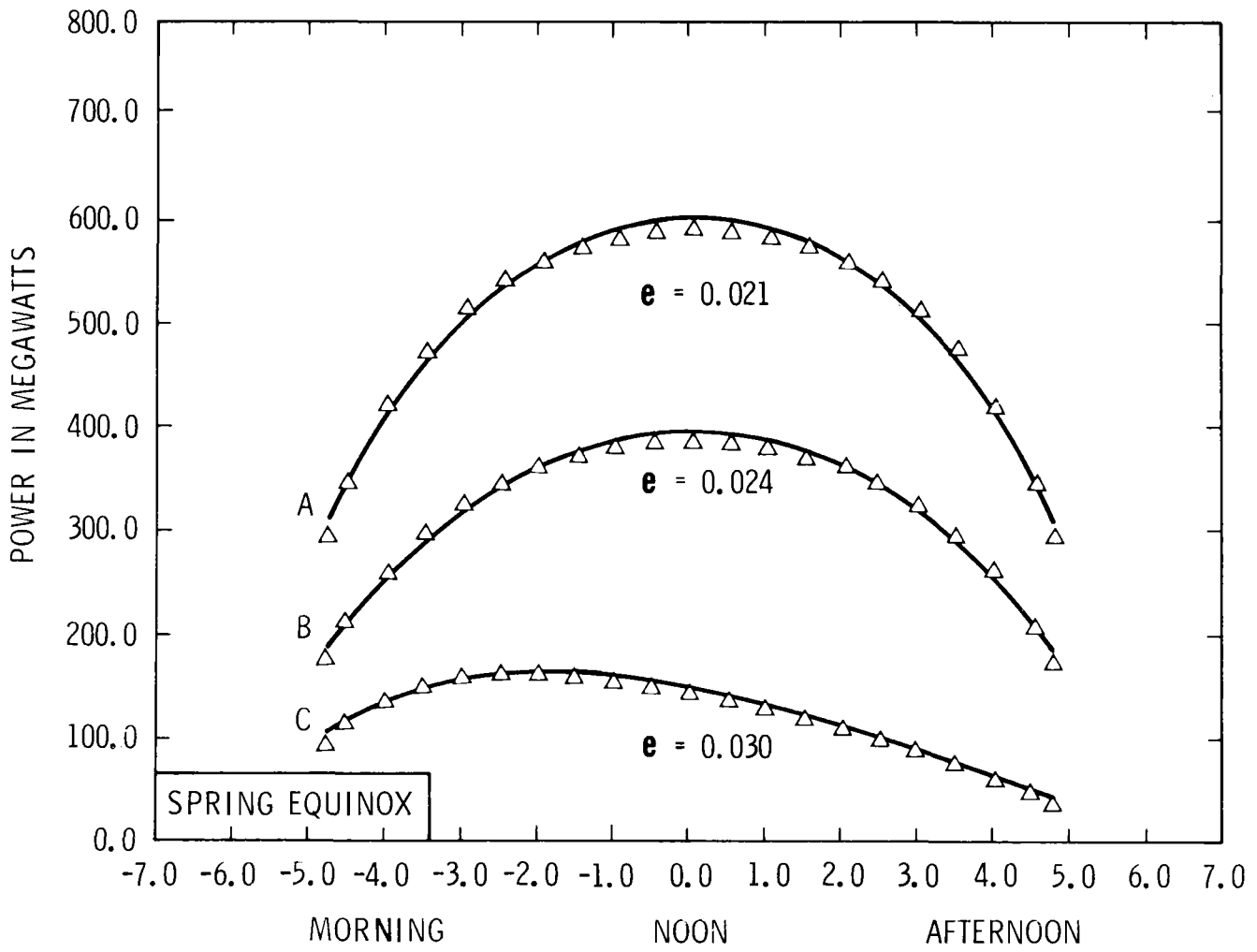


Figure 6. Same as Figure 5 but for March 21.

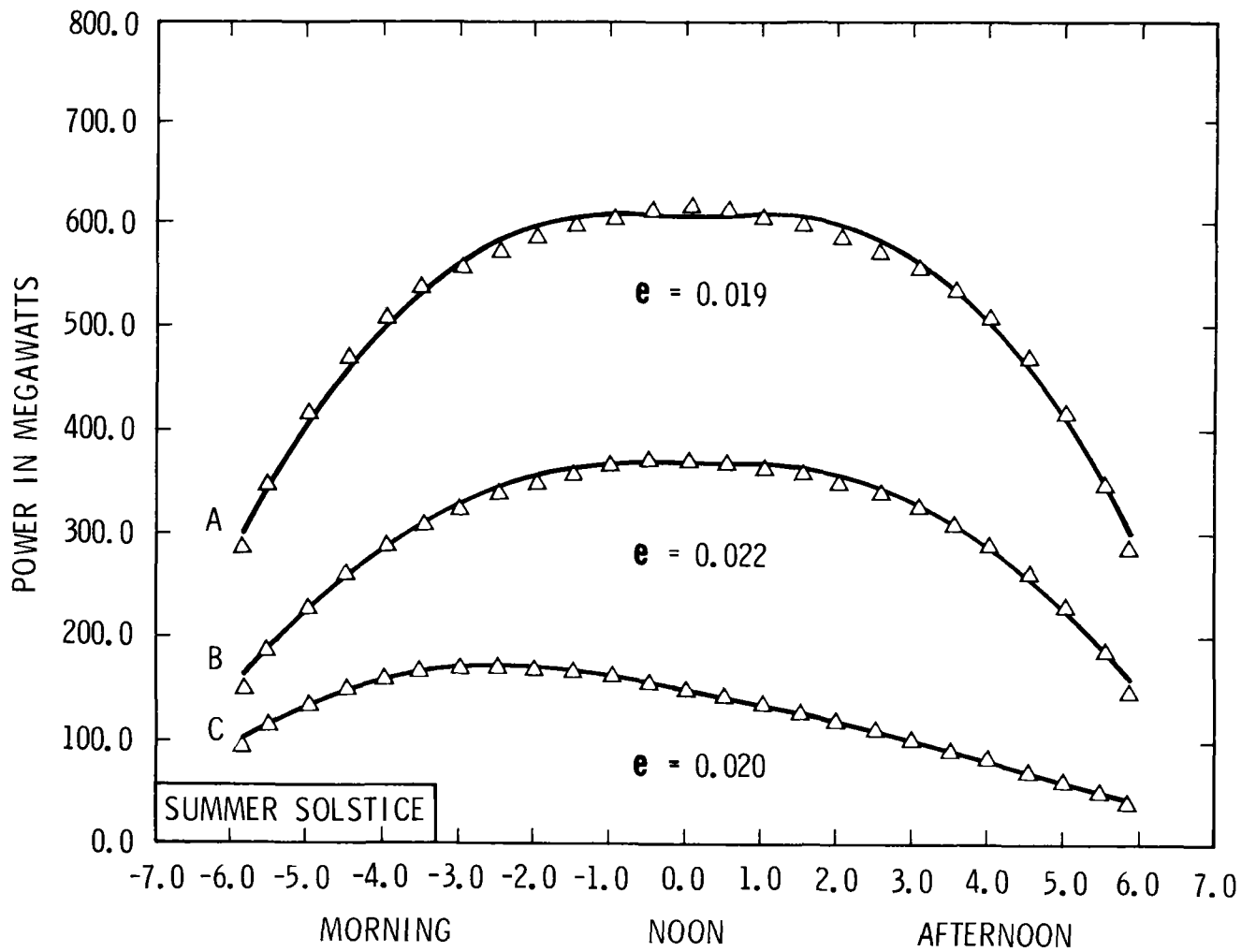


Figure 7. Same as Figure 5 but for June 22.

As a more complete demonstration of the capabilities of the SCRAM code, we have calculated the daily performance of a particular heliostat field using TMY weather data for Albuquerque, NM, and the program outlined in Figure 4. The field layout and results are shown in Figure 8. Total mirror area was 75,000 m², and the tower height (calculated by the program) was 62.5 m. The annual average efficiency for the field (defined by the dashed line in Figure 7) was 52%. If receiver losses are not included, this number becomes 65%.

This calculation required several thousand calls to the SCRAM subroutine (once every daylight hour for one year), and took 9.3 seconds of CDC 6600 computer time. Though it is difficult to make direct comparisons, this is 1,000-3,000 faster than DELSOL, which in turn is probably faster than other existing PFCR codes.⁹ (It should also be noted that most of this 9.3 seconds is occupied by reading the weather data; only 3 seconds were actually used by the simulation programs.)

VI. Conclusion

The fast computer model described here is not intended to replace the large simulation codes for all applications. By increasing the speed of the calculation, we have limited the flexibility of the analysis as well. For example, we cannot change the geometry of the heliostats without generating a new data base. (See Appendix A for a discussion of which parameters can be changed and which cannot.)

However, the speed of the SCRAM model makes certain types of analysis possible which would require prohibitive amounts of time with the simulation codes. For example, optimizing the tower location on a fixed piece of land based on annual performance with realistic weather data would require more than 10⁵ evaluations of field performance, feasible for SCRAM, but not for any other existing code. Another suitable application would be the mutual optimization of field boundaries and storage size when a realistic load is considered and the effects of transients and local weather are treated.

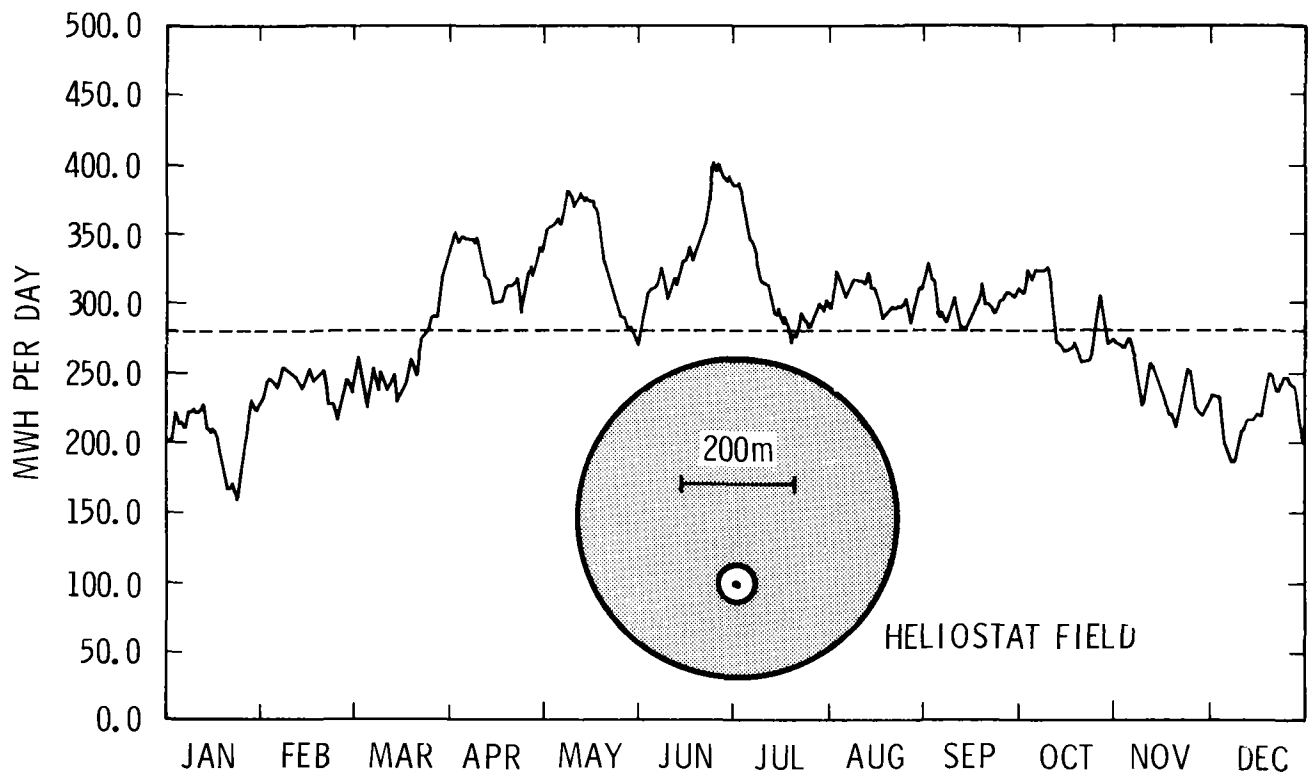


Figure 8. Annual performance of a PFCR system based on TMY (Typical Meteorological Year) data for Albuquerque, NM. Receiver losses are included. The mirror area is 75,000 m² and the tower height (calculated by the program) is 62.5 m.

Another advantage is the simplicity of the computation. Since there is no need for a great deal of memory or speed, the model can be implemented on any size computer. Another possibility is that, once the field has been specified and the values of X_{ij} have been determined by numerical integration, the calculation of field response can be performed on an even smaller machine (e.g., microcomputer or calculator), since only multiplication and addition is required. As central receiver systems evolve in stages from concept to prototype to potentially significant penetration of the electrical generation and industrial process heat markets, we can expect the importance of these and many other applications of a fast computer model to increase.

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APPENDIX A

Heliostat Characteristics and
Scaling Relationships

DELSOL requires a large number of input parameters to specify the geometry and characteristics of the heliostats. The values for the data which generated the SCRAM model presented in Table III are the default values defined in SAND79-8215 (Reference 3). The basic configuration is a rectangular structure 7.4 m x 7.4 m with focused, canted mirrors covering 90% of this area. The tower height is defined to be the distance from the horizontal rotation axis so that the height of the heliostat above the ground is not needed. Mirror reflectivity is .89.

The field layout is basically a radial stagger pattern, also described in Reference 3. The heliostat density (mirror area/land area) as a function of field position is given by

$$\rho = \frac{D_M \cdot (7.4)^2}{\Delta R \Delta A / 2} \quad (A1)$$

where D_M is the ratio of mirror area to heliostat area, and ΔR and ΔA are the radial and azimuthal separations (in meters) respectively. These separations are given by the formulas

$$\Delta R = (1.144 \cot \theta_L - 1.094 + 3.068 \theta_L - 1.1256 \theta_L^2) \cdot 7.4 \quad (A2)$$

$$\Delta A = (1.749 + 0.6396 \theta_L + \frac{0.02873}{\theta_L - 0.04902}) \cdot 7.4 \quad (A3)$$

where $\theta_L = \cot^{-1}(r)$ with r in tower heights. These formulas are based on studies of optimum heliostat layout by workers at University of Houston.¹⁰ They are incorporated into the subroutine DENS in Appendix C.

The formula used for atmospheric attenuation is Eq. III. E-1 from Reference 3 corresponding to a clear day in Barstow, California. This formula, and others which could be substituted for different conditions and different locations can be found in Reference 11.

Most of the parameters characterizing the heliostats and layout are not variable once the SCRAM model, C_{ijk} , has been calculated. However, there

are a few of them which can be varied without requiring a new model calculation because of the existence of simple scaling relations. Below are listed the parameters in question and the procedure required to take advantage of the scaling relation.

- (1) Tower height: incorporated in existing algorithm (see Appendix D).
- (2) Mirror reflectivity: multiply each C_{ijk} by $(M_r/.89)$ where M_r is mirror reflectivity.
- (3) Different atmospheric attenuation models: C_{ijk} unchanged; subroutine DENS is modified.
- (4) D_M , ratio of mirror area to heliostat area: C_{ijk} unchanged; parameter DENSMIR = D_M in subroutine DENS is changed from $D_M = .897$.
- (5) Receiver losses (spillage, reflection, radiation, convection and conduction): incorporate into separate receiver model instead of the simplistic one used in SCRAM. (See Appendix C.)

We begin with Eq. (7) for the local field efficiency:

$$\eta_p = (r, \xi, \theta) = \sum_{i, j, k=1, 4} C_{ijk} V_i(\theta) U_j(\xi) W_k(r) \quad (B1)$$

Integrating this over the field area as in Eq. (6) gives

$$C_f(\theta, \phi) = \sum_{i, k=1, 4} \int d\beta \int r dr W_k(r) V_i(\theta) \eta_{ik}(\phi, \beta) \quad (B2)$$

where we have introduced

$$\eta_{ik}(\phi, \beta) = \sum_{j=1, 4} C_{ijk} \{ \cos((j-1)\phi) \cos((j-1)\beta) + \sin((j-1)\phi) \sin((j-1)\beta) \} \quad (B3)$$

The definitions in Table II and the trigonometric expansion for $\cos(\phi - \beta)$ have been used to obtain Eq. (B3). We now restrict the field boundaries to two closed curves--an inner radius $r_1(\beta)$ and an outer radius $r_2(\beta)$ so that we can define X_{ij} as follows:

$$X_{i1} = \int_{r_1}^{r_2} r dr \sum_{k=1, 4} r W_k(r) C_{i1k} \quad (B4)$$

$$X_{ij+1} = \int d\beta \sum_{k=1, 4} \int_{r_1}^{r_2} r W_k(r) dr C_{ij+1k} \cos(j\beta) \quad (B5)$$

$$X_{ij+4} = \int d\beta \sum_{k=1, 4} \int_{r_1}^{r_2} r W_k(r) dr C_{ij+1k} \sin(j\beta) \quad (B6)$$

where $j = 1, 2, 3$ in the last two equations. This allows us to write (B2) in the form

$$C_f(\theta, \phi) = \sum_{\substack{i=1,4 \\ j=1,7}} X_{ij} V_i(\theta) U_j(\phi) \quad (B7)$$

which is Eq. (8) from the text. Since X_{ij} does not depend on θ or ϕ , it needs to be calculated only once for each field.

C.1 Thermal and Reflection Losses

As discussed in the text, the principal reason for separating field losses from receiver losses in this analysis is that the nature of the physics and the accuracy of existing calculational tools for the processes involved is significantly different for the two categories. For example, the best current models for calculating convective heat loss are not nearly as well developed as our ability to do classical optics calculations. (However, it is to be expected that in the next few years rapid progress in understanding the thermal properties of PFCR systems will result from the DOE research and demonstration program.)

In this section, a simplified model for receiver performance is described, partly as a demonstration of how such a model would be incorporated and partly to provide a more complete PFCR simulation for applications which do not require high levels of accuracy for receiver performance.

We have identified receiver losses as consisting of five terms, as indicated in Table I. In our receiver model, the loss due to reflection is treated simply by assigning a value to η_7 , identified as ROPT in SCRAM. This is suitable for both external and cavity receivers, though for the latter it is the effective absorptivity of the cavity (before re-radiation) which is being modeled.

The three thermal terms, η_8 , η_9 , and η_{10} , are also dealt with in a very simple manner. It is assumed that the receiver is held at a fixed temperature regardless of insulation level (by controlling the flow of heat transfer fluid) so that the power loss is simply a constant in time. The quantity ROTHER in subroutine SCRAM specifies this loss in units of power consistent with the insulation units. No attempt has been made to model the dependence of this quantity on temperature, wind speed, etc. Thus the receiver model for the terms $\eta_7 \cdot \eta_8 \cdot \eta_9 \cdot \eta_{10}$ is:

$$\text{Collected Power} = \text{Incident Power} \cdot \text{ROPT} - \text{ROTHER} \quad (\text{C1})$$

For the example of Figure 8, the quantities ROTHER = 2 MW and ROPT = .87 were used. The former was based on an arbitrary assumption of a thermal

loss equal to 5% of the peak collected power, while the latter is a rather conservative estimate of absorptivity for an external receiver.

C.2 Spillage

It is natural to treat η_{7-10} as attributes of the receiver, but it is somewhat arbitrary to consider η_6 , the spillage term, a property of the receiver. It is actually a property of the interface between the field and the receiver, and as such it presents some problems for a fast approximation procedure.

Strictly speaking, spillage is an optical process and DELSOL's treatment of it should have the same high level of accuracy as the other optical terms in the field efficiency. However, there are important advantages to including spillage in the receiver model and treating it in an approximate way.

Firstly, it should be noted that spillage depends on a large number of parameters, including receiver dimensions and heliostat errors, as well as r , θ , and ξ . In fact, for a typical, simple system (external cylindrical receiver, canted and focused heliostats) DELSOL requires fifteen input parameters for the heliostats and three for the receiver. Of these eighteen parameters, fully fourteen affect only the spillage losses. Thus, if the polynomial η_p does not include spillage, there is a substantial increase in the generality of the polynomial.

Secondly, spillage in a well-designed system typically makes a very small contribution to PFCR losses. The reason is that the heliostat field represents such a large proportion of system costs that it is cost-effective to design receivers which intercept almost all of the reflected light. For example, in a typical set of DELSOL-optimized PFCR systems, the annual spillage loss is less than 1%.

On the other hand, we do not want to disregard spillage in certain types of problems (e.g. field boundary optimization). The procedure to be adopted here is based on the simple assumption of ignoring the dependence of spillage on θ and ξ (which is known to be a weak dependence for

typical systems). With spillage a function only of r , it can be treated in the same manner as η_5 , atmospheric attenuation (included in the subroutine FLDSET).

The function used for our optional receiver model is based on DELSOL results using canted, focused heliostats and an external cylindrical receiver. As before, the parameters used are the default values of Reference 3. The functional form is

$$\begin{aligned} \eta_7 &= 1 && (r \leq 4.2) \\ \eta_7 &= 1 - (r - 4.2)^{1.89} && (r > 4.2) \end{aligned} \tag{C2}$$

An analysis of DELSOL runs indicates that errors introduced by this approximate treatment would typically be less than 10% of the spillage, hence less than .1% of the energy collected.

One important advantage of this procedure is a substantial increase in flexibility. Thus, a different receiver geometry or different heliostat characteristics would require a new spillage function, but not a new field polynomial. Furthermore, the spillage calculation only requires the analysis of a single heliostat located at varying values of r (and preferably at $\xi = 90^\circ$).

If single-cavity receivers or receivers whose parameters are very far from optimum are used, this treatment of spillage may not be adequate. For example, a south-facing cavity will obviously not accept light from the north part of the field. It might then be possible to introduce an angle-dependent spillage factor (which would be included in FLDSET). Similarly, receiver design and performance studies will probably require flux maps on the receiver surface and the SCRAM approach would be unsuitable. However, for most system studies and field design studies involving receivers designed for surround fields, the treatment of spillage outlined above should be entirely adequate.

APPENDIX D

Computer Subroutines for Implementing the SCRAM Model

The following subroutines were used with different executive programs to generate the data for Figures 5-8. The flow chart for the program generating Fig. 8 is shown in Fig 4. Two library subroutines for numerical quadrature were used: QNC3 for function-generated integrands, and AVINT for integrating discrete data sets. They should be replaced by appropriate substitutes for systems other than that at Sandia Laboratories.

```

SUBROUTINE TUNT (RI,RO,ANG,JMX,AM)
C
C TUNT CALCULATES THE TOWER HEIGHT, GIVEN THE FIELD BOUNDARIES
C (NORMALIZED TO THT) AND THE REQUIRED MIRROR AREA.
C THE FUNCTION DENS IS CALLED TO SPECIFY THE MIRROR DENSITY AS A
C FUNCTION OF RADIUS.
C
EXTERNAL DENS
COMMON /DEN/ XRAD,IFLAG,THT
DIMENSION H(50), RI(50), RO(50), ANG(50)
XRAD=2.0
ERR=.001
IFLAG=1
DO 10 K=1,JMX
10 CALL QNC3 (DENS,RI(K),RO(K),ERR,H(K),IERR)
CALL AVINT (ANG,H,JMX,ANG(1),ANG(JMX),AMNORM,IER2)
THT=(AM/AMNORM)**.5
PRINT 20, AM,THT
RETURN
20 FORMAT (/13H MIRROR AREA=,E10.3,10H 50 METERS,16H --TOWER HEIGHT=,
1F7.3,7H METERS)
END
C
C
C
SUBROUTINE FLDSET (RI,RO,ANG,JMX,X)
C
C FLDSET USES THE GENERAL MODEL C(I,J) AND THE FIELD SPECIFICATION
C DEFINED BY RI, RO, ANG, AND THT TO
C CALCULATE A FIELD MODEL X(I,J) WHICH IS USED BY SCRAM FOR
C PERFORMANCE CALCULATIONS.
C
EXTERNAL DENS
COMMON /DEN/ XRAD,IFLAG,THT
DIMENSION RI(50), RO(50), ANG(50), H1(50), H2(50), CJO(25), JO(25)
1, X(4,7), ZI(4,50), U(6), C(4,4,4)
DATA CJO/-.92596,.546639,.9343,.11135,.56558,.082405,-.02308,.0742
164,-.005781/
DATA JO/9,2,3,5,9,11,40,41,44,49/
PRINT 90
PRINT 70
IFLAG=0
XRAD=2.
ERR=.001
IO=1
ITERM=2
C-----RECONSTRUCT THREE INDEX CODE FROM CJO(CO)---
DO 10 KRANK=3,12
DO 10 I=1,4
DO 10 J=1,4
DO 10 K=1,4
IF (I+J+K.NE.KRANK) GO TO 10
IO=IO+1
C(I,J,K)=0.
IF (IO.NE.JO(ITERM)) GO TO 10
C(I,J,K)=CJO(ITERM-1)
PRINT 80, JO(ITERM),I,J,K,C(I,J,K)

```

```

      ITERM=ITERM+1
10  CONTINUE
C  -----PERFORM R INTEGRALS-----
      DO 30 K=1,4
      DO 20 KK=1,JMX
      XRAD=K
      CALL QNC3 (DENS,RI(KK),RO(KK),ERR,ANS,IERR)
20  ZI(K,KK)=ANS
30  CONTINUE
      DO 60 I=1,4
      DO 60 J=1,4
      X(I,J)=0.
      ZJ=J-1
      DO 50 KK=1,JMX
      DH=0.
      DO 40 K=1,4
40  DH=DH+C(I,J,K)*ZI(K,KK)
      H1(KK)=DH*COS(ZJ*ANG(KK))
50  H2(KK)=DH*SIN(ZJ*ANG(KK))
      CALL AVINT (ANG,H1,JMX,ANG(1),ANG(JMX),PO1,IER1)
      CALL AVINT (ANG,H2,JMX,ANG(1),ANG(JMX),PO2,IER2)
      X(I,J)=PO1*THT**2
60  IF (J.NE.1) X(I,J+3)=PO2*THT**2
      RETURN
C
70  FORMAT (/24H JO I J K CO/)
80  FORMAT (4I4,E14.7)
90  FORMAT (/22H SCRAM MODEL C(I,J,K):/)
      END
C
C
C  SUBROUTINE SCRAM (SAA,SPA,AM,DN,PWR,X)
C
C  SCRAM USES THE FIELD MODEL X(I,J) TO CALCULATE THE ENERGY
C  INCIDENT ON THE RECEIVER. IT CAN ALSO (IF IREC=1)
C  CALCULATE POWER LOST AT THE RECEIVER AND RETURNS NET POWER
C  COLLECTED TO ANYFLD.
C
      COMMON /REC/ ROPT,RTHET,IREC
      DIMENSION E(4), X(4,7)
      E(1)=1.
      E(2)=COS(SPA)
      E(3)=SIN(SPA)
      E(4)=E(3)*E(2)
      ETA=0.
      DO 30 I=1,4
      EF=X(I,1)
      DO 20 J=1,3
      FJ=J
      IF (X(I,J+1).EQ.0.) GO TO 10
      EF=X(I,J+1)*COS(FJ*SAA)+EF
10  IF (X(I,J+4).EQ.0.) GO TO 20
      EF=EF+X(I,J+4)*SIN(FJ*SAA)
20  CONTINUE
30  ETA=ETA+EF*E(I)
      PWR=ETA*DN
C  -----OPTIONAL RECEIVER MODEL---

```

```

IF (IREC.EQ.0) RETURN
PUR=PW*ROPT-R*THEM
IF (PUR.LT.0.) PUR=0.
RETURN
END

C
C
C
FUNCTION DENS(RAD)
C
C   DENS CALCULATES MIRROR DENSITY AS A FUNCTION OF RADIUS USING
C   UNIVERSITY OF HOUSTON FORMULAE FOR OPTIMUM
C   HELIOSTAT LAYOUT.  IT ALSO CALCULATES ATMOSPHERIC ATTENUATION
C   (WHEN IFLAG=0.)
C
COMMON /DEN/ XRAD,IFLAG,THT/REC/ROPT,R*THEM,IREC
-----IFLAG = 1 EXCLUDES ATMOSPHERIC ATTENUATION-----
DENSMIR=.89
THETA=ATAN(1.0211/RAD)
RSEP=1.14424/TAN(THETA)-1.09352+3.06836*THETA-1.12556*THETA**2
ASEP=1.74909+.63964*THETA+.02873/(THETA-.04902)
PDR=((2.0*DENSMIR)/(RSEP*ASEP))
IF (IFLAG.EQ.1) GO TO 10
RANGE=RAD*THT/1000.0
ATMOS=.6789+10.46*RANGE-1.70*RANGE**2+.2845*RANGE**3
PDR=PDR*(1.0-ATMOS/100.0)
-----SPILLAGE MODEL-----
10 IF (IREC.EQ.1.AND.RAD.GT.4.2) PDR=PDR*(1.-.01*(RAD-4.2)**1.89)
DENS=PDR/RAD
DENS=DENS*RAD**XRAD
RETURN
END

C
C
C

```

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